

Starting from the usual force equation:

$$\ddot{r} = -G\frac{4}{3}\pi\rho r \quad (1)$$

the energy density of blackbody light:

$$\rho_\gamma = u/c^2 = \frac{4\sigma T^4}{c^3} \quad (2)$$

the usual scaling: $r = uR$ and how the light's temperature is affected by expansion:

$$T \longrightarrow \frac{T_0}{R} \quad (3)$$

we have:

$$\ddot{R} = -\frac{G\frac{4}{3}\pi\rho_\gamma}{R^3} \quad (4)$$

$$\frac{1}{2}\dot{R}^2 = \frac{G\frac{4}{3}\pi\rho_\gamma}{2R^2} \quad (5)$$

where an integration constant has been taken as zero. If we include radiation pressure, this result is doubled. Combining this result with previous:

$$\dot{R}^2 = \frac{G\frac{8}{3}\pi\rho_\gamma}{R^2} + \frac{G\frac{8}{3}\pi\rho_m}{R} + \frac{1}{3}\Lambda c^2 R^2 \quad (6)$$

where ρ_m and ρ_γ are the NOW ($R = 1$) values of mass and radiation density. Using the usual definitions

$$H_0 = 67.9 \text{ km/s/Mpc} \quad (7)$$

$$T_0 = 2.725 \text{ K} \quad (8)$$

$$\Omega_\gamma = \frac{8\pi G\rho_\gamma}{3H_0^2} = \frac{8\pi G4\sigma T_0^4}{3c^3 H_0^2} = 5.4 \times 10^{-4} \quad \text{adding neutrinos} \rightarrow 8.4 \times 10^{-4} \quad (9)$$

$$\Omega_m = \frac{8\pi G\rho_m}{3H_0^2} = .306 \quad (10)$$

$$\Omega_\Lambda = \frac{\Lambda c^2}{3H_0^2} = .694 \quad (11)$$

we get the general differential equation

$$\dot{R}^2 = H_0^2 \left(\frac{\Omega_\gamma}{R^2} + \frac{\Omega_m}{R} + \Omega_\Lambda R^2 \right) \quad (12)$$

In the early universe (say $R < .01$) we neglect the last term. Separating the resulting differential

equation

$$\int_0^R \frac{dR}{\sqrt{\frac{\Omega_\gamma}{R^2} + \frac{\Omega_m}{R}}} = H_0 t \quad (13)$$

$$\int_0^R \frac{R dR}{\sqrt{\Omega_\gamma + \Omega_m R}} = H_0 t \quad (14)$$

$$\frac{2}{\Omega_m^2} \left[\frac{1}{3} X^{3/2} - \Omega_\gamma X^{1/2} \right]_0^R = H_0 t \quad (15)$$

$$\frac{2}{\Omega_m^2} \left\{ \frac{1}{3} \left(\sqrt{\Omega_\gamma + \Omega_m R} \right)^3 - \Omega_\gamma \sqrt{\Omega_\gamma + \Omega_m R} + \frac{2}{3} \Omega_\gamma^{3/2} \right\} = H_0 t \quad (16)$$

$$\frac{R^2}{2\sqrt{\Omega_\gamma}} - \frac{\Omega_m R^3}{6\Omega_\gamma^{3/2}} + \mathcal{O}(R^4) = H_0 t \quad (17)$$

In the last line the result has been series expanded to clearly show for early times

$$R \approx \sqrt{2\sqrt{\Omega_\gamma} H_0 t} \propto t^{1/2} \quad (18)$$

```
Integrate[r/Sqrt[OmegaG+OmegaM r ],{r,0,R},GenerateConditions->False]
Series[%,{R,0,3}]
```

$$\text{Out [4]} = \frac{R^2}{2 \text{Sqrt}[\text{OmegaG}]} - \frac{\text{OmegaM} R^3}{6 \text{OmegaG}^{3/2}} + \mathcal{O}[R]$$

```
out=%%/H0 /. {OmegaM->.306,H0->1/14.4,OmegaG->.000084}
FindRoot[out==.00038,{R,.001}]
```

The last line allows finding R when $t = .00038$ (i.e., 380,000 years). Find z and $T = T_0/R$ at that time. This should be an improvement over the last homework result. Note from the choice of units for H_0 we are using time in billions of years.

Matter and light carry equal weight when

$$\frac{\Omega_\gamma}{R^2} = \frac{\Omega_m}{R} \quad (19)$$

When is that?

$t = 3$ minutes is the time of nuclear synthesis; find R and T at that time.